This example is set up in Antykwa Półtawskiego. It uses:
e[0T4,0T1]\{fontenc\}\usepackage\{txfonts\}\usepackage[upright]\{txgreeks\}\usepackage\{antpolt\}\usepackage[defaultmathsizes,nolessnomore]\{mathastext\}Typesetwithmathastext1.13b(2011/03/15).undefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefinedundefined

To illustrate some Hilbert Space properties of the co-Poisson summation, we will assume $K=\mathbf{Q}$. The components ( $\mathrm{a}_{\mathrm{v}}$ ) of an adele a are written $\mathrm{a}_{\mathrm{p}}$ at finite places and $\mathrm{a}_{\mathrm{r}}$ at the real place. We have an embedding of the Schwartz space of test-functions on $\mathbf{R}$ into the Bruhat-Schwartz space on $\mathbf{A}$ which sends $\psi(x)$ to $\varphi(a)=\prod_{p} \mathbf{1}_{\left|a_{p}\right|_{p} \leq 1}\left(a_{p}\right) \cdot \psi\left(a_{r}\right)$, and we write $E_{\mathbf{R}}^{\prime}(\mathrm{g})$ for the distribution on $\mathbf{R}$ thus obtained from $\mathrm{E}^{\prime}(\mathrm{g})$ on $\mathbf{A}$.

Theorem 1. Let g be a compact Bruhat-Schwartz function on the ideles of Q. The co-Poisson summation $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ is a square-integrable function (with respect to the Lebesgue measure). The $\mathrm{L}^{2}(\mathbf{R})$ function $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ is equal to the constant $-\int_{\mathbf{A}^{\times}} \mathrm{g}(\mathrm{v})|\mathrm{v}|^{-1 / 2} \mathrm{~d}{ }^{*} \mathrm{v}$ in a neighborhood of the origin.
Proof. We may first, without changing anything to $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$, replace g with its average under the action of the finite unit ideles, so that it may be assumed invariant. Any such compact invariant $g$ is a finite linear combination of suitable multiplicative translates of functions of the type $g(v)=\prod_{p} \mathbf{1}_{\mid \mathrm{v}_{\mathrm{p} \mid \mathrm{p}}=1}\left(\mathrm{v}_{\mathrm{p}}\right) \cdot \mathrm{f}\left(\mathrm{v}_{\mathrm{r}}\right)$ with $f(t)$ a smooth compactly supported function on $\mathbf{R}^{\times}$, so that we may assume that g has this form. We claim that:

$$
\int_{\mathbf{A}^{\times}}|\varphi(\mathrm{v})| \sum_{\mathrm{q} \in \mathbf{Q}^{\times}}|\mathrm{g}(\mathrm{qv})| \sqrt{|\mathrm{v}|} \mathrm{d}^{*} \mathrm{v}<\infty
$$

Indeed $\sum_{\mathbf{q}^{\prime} \in \mathbf{Q}^{\times}}|\mathrm{g}(\mathrm{qv})|=|\mathrm{f}(|\mathrm{v}|)|+|\mathrm{f}(-\mid \mathrm{v})|$ is bounded above by a multiple of $|\mathrm{v}|$. And $\int_{\mathbf{A}^{\times}}|\varphi(\mathrm{v}) \| \mathrm{v}|^{3 / 2} \mathrm{~d}^{*} \mathrm{v}<\infty$ for each Bruhat-Schwartz function on the adeles (basically, from $\left.\prod_{p}\left(1-p^{-3 / 2}\right)^{-1}<\infty\right)$. So

$$
\begin{aligned}
& E^{\prime}(g)(\varphi)=\sum_{q \in \mathbf{Q}^{\times}} \int_{\mathbf{A}^{\times}} \varphi(v) g(q) \sqrt{|v|} d^{*} v-\int_{\mathbf{A}^{\times}} \frac{g(v)}{\sqrt{|v|}} d^{*} v \int_{\mathbf{A}} \varphi(x) d x \\
& E^{\prime}(g)(\varphi)=\sum_{q \in \mathbf{Q}^{\times}} \int_{\mathbf{A}^{\times}} \varphi(v / q) g(v) \sqrt{|v|} d^{*} v-\int_{\mathbf{A}^{\times}} \frac{g(v)}{\sqrt{|v|}} d^{*} v \int_{\mathbf{A}} \varphi(x) d x
\end{aligned}
$$

Let us now specialize to $\varphi(\mathrm{a})=\prod_{\mathrm{p}} \mathbf{1}_{\mid \mathrm{ap}_{\mathrm{p}} \leq 1}\left(\mathrm{a}_{\mathrm{p}}\right) \cdot \psi\left(\mathrm{a}_{\mathrm{r}}\right)$. Each integral can be evaluated as an infinite product. The finite places contribute 0 or 1 according to whether $\mathrm{q} \in \mathbf{Q}^{\times}$satisfies $|\mathrm{q}|_{\mathrm{p}}<1$ or not. So only the inverse integers $\mathrm{q}=1 / \mathrm{n}, \mathrm{n} \in \mathbf{Z}$, contribute:

$$
\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})(\psi)=\sum_{\mathrm{n} \in \mathbf{Z}^{\times}} \int_{\mathbf{R}^{\times}} \psi(\mathrm{nt}) \mathrm{f}(\mathrm{t}) \sqrt{|\mathrm{tt}|} \frac{\mathrm{dt}}{2|t|}-\int_{\mathbf{R}^{\times}} \frac{\mathrm{f}(\mathrm{t})}{\sqrt{|t|}} \frac{\mathrm{dt}}{2|t|} \int_{\mathbf{R}} \psi(\mathrm{x}) \mathrm{dx}
$$

We can now revert the steps, but this time on $\mathbf{R}^{\times}$and we get:

$$
E_{\mathbf{R}}^{\prime}(\mathrm{g})(\psi)=\int_{\mathbf{R}^{\times}} \psi(\mathrm{t}) \sum_{\mathrm{n} \in \mathbf{Z}^{\times}} \frac{\mathrm{f}(\mathrm{t} / \mathrm{n})}{\sqrt{|\mathrm{n}|}} \frac{\mathrm{dt}}{2 \sqrt{|t|}}-\int_{\mathbf{R}^{\times}} \frac{\mathrm{f}(\mathrm{t})}{\sqrt{|t|}} \frac{\mathrm{dt}}{2|t|} \int_{\mathbf{R}} \psi(\mathrm{x}) \mathrm{dx}
$$

Let us express this in terms of $\alpha(\mathrm{y})=(\mathrm{f}(\mathrm{y})+\mathrm{f}(-\mathrm{y})) / 2 \sqrt{|\mathrm{y}|}$ :

$$
E_{\mathbf{R}}^{\prime}(\mathrm{g})(\psi)=\int_{\mathbf{R}} \psi(\mathrm{y}) \sum_{\mathrm{n} \geq 1} \frac{\alpha(\mathrm{y} / \mathrm{n})}{\mathrm{n}} \mathrm{dy}-\int_{0}^{\infty} \frac{\alpha(\mathrm{y})}{\mathrm{y}} \mathrm{dy} \int_{\mathbf{R}} \psi(\mathrm{x}) \mathrm{dx}
$$

So the distribution $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ is in fact the even smooth function

$$
E_{R}^{\prime}(\mathrm{g})(\mathrm{y})=\sum_{\mathrm{n} \geq 1} \frac{\alpha(\mathrm{y} / \mathrm{n})}{\mathrm{n}}-\int_{0}^{\infty} \frac{\alpha(\mathrm{y})}{\mathrm{y}} \mathrm{dy}
$$

As $\alpha(\mathrm{y})$ has compact support in $\mathbf{R} \backslash\{0\}$, the summation over $\mathrm{n} \geq 1$ contains only vanishing terms for $|y|$ small enough. So $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ is equal to the con-
 To prove that it is $L^{2}$, let $\beta(y)$ be the smooth compactly supported function $\alpha(1 / y) / 2|y|$ of $y \in \mathbf{R}(\beta(0)=0)$. Then $(y \neq 0)$ :

$$
\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})(\mathrm{y})=\sum_{\mathrm{n} \in \mathbf{Z}} \frac{1}{|\mathrm{y}|} \beta\left(\frac{\mathrm{n}}{\mathrm{n}}\right)-\int_{\mathbf{R}} \beta(\mathrm{y}) \mathrm{dy}
$$

From the usual Poisson summation formula, this is also:

$$
\sum_{n \in \mathbf{Z}} \gamma(n y)-\int_{\mathbf{R}} \beta(y) d y=\sum_{n \neq 0} \gamma(n y)
$$

where $\gamma(\mathrm{y})=\int_{\mathbf{R}} \exp (\mathrm{i} 2 \pi y \mathrm{w}) \beta(\mathrm{w}) \mathrm{d} \boldsymbol{w}$ is a Schwartz rapidly decreasing function. From this formula we deduce easily that $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})(\mathrm{y})$ is itself in the Schwartz class of rapidly decreasing functions, and in particular it is is square-integrable.

It is useful to recapitulate some of the results arising in this proof:
Theorem 2. Let g be a compact Bruhat-Schwartz function on the ideles of $\mathbf{Q}$. The co-Poisson summation $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ is an even function on $\mathbf{R}$ in the Schwartz class of rapidly decreasing functions. It is constant, as well as its Fourier Transform, in a neighborhood of the origin. It may be written as

$$
E_{R}^{\prime}(\mathrm{g})(\mathrm{y})=\sum_{\mathrm{n} \geq 1} \frac{\alpha(\mathrm{y} / \mathrm{n})}{\mathrm{n}}-\int_{0}^{\infty} \frac{\alpha(\mathrm{y})}{\mathrm{y}} \mathrm{dy}
$$

with a function $\alpha(y)$ smooth with compact support away from the origin, and conversely each such formula corresponds to the co-Poisson summation $\mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})$ of a compact Bruhat-Schwartz function on the ideles of $\mathbf{Q}$. The Fourier transform $\int_{\mathbf{R}} \mathrm{E}_{\mathbf{R}}^{\prime}(\mathrm{g})(\mathrm{y}) \exp (\mathrm{i} 2 \pi \mathrm{wy})$ dy corresponds in the formula above to the replacement $\alpha(\mathrm{y}) \mapsto \alpha(1 / \mathrm{y}) /|\mathrm{y}|$.

Everything has been obtained previously.

